

# The Necker-Zeno Model

Harald Atmanspacher, IGPP Freiburg

Collaboration with T. Filk, J. Kornmeier, H. Römer

- 1 Introduction
- 2 Necker-Zeno Model for Bistable Perception
- 3 Empirical Confirmation
- 4 Temporal Nonlocality
- 5 Selected References

## Some Remarks

- psychology is different from neuroscience

## Some Remarks

- psychology is different from neuroscience
- mathematics is more than data processing

## Some Remarks

- psychology is different from neuroscience
- mathematics is more than data processing
- mathematical precision is more than quantitative

## Some Remarks

- psychology is different from neuroscience
- mathematics is more than data processing
- mathematical precision is more than quantitative

Mathematics serves the precise formulation of conceptual questions in terms of **abstract structures** (algebras, graphs, etc.).

Data processing includes the **numerical quantification** of observables, statistical analysis of measurement results, etc.

Observational processes are **interactions** of an observing system  $O$  with an observed system  $S$  (state  $\psi$ , observables  $A, B, \dots$ ):

- (i) weak interaction: no significant effect of  $O$  on  $S$ ,
- (ii) strong interaction: effect of  $O$  on  $S$  makes a difference.

Observational processes are **interactions** of an observing system  $O$  with an observed system  $S$  (state  $\psi$ , observables  $A, B, \dots$ ):

- (i) weak interaction: no significant effect of  $O$  on  $S$ ,
- (ii) strong interaction: effect of  $O$  on  $S$  makes a difference.

Physics:

- (i) classical case,  $AB\psi = BA\psi$  **commutative**
- (ii) quantum case,  $AB\psi \neq BA\psi$  **non-commutative**



Observational processes are **interactions** of an observing system  $O$  with an observed system  $S$  (state  $\psi$ , observables  $A, B, \dots$ ):

- (i) weak interaction: no significant effect of  $O$  on  $S$ ,
- (ii) strong interaction: effect of  $O$  on  $S$  makes a difference.

Physics:

- (i) classical case,  $AB\psi = BA\psi$  **commutative**
- (ii) quantum case,  $AB\psi \neq BA\psi$  **non-commutative**

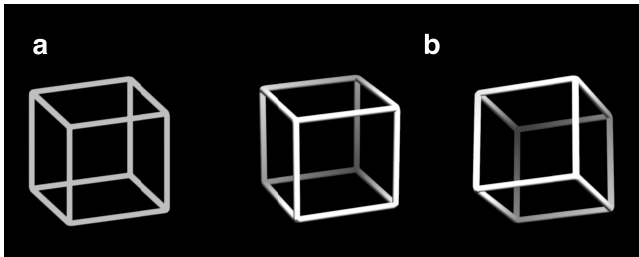
Psychology:

Almost every action of  $O$  entails a significant effect on  $S$ .  
Non-commutativity is the rule rather than the exception.

→ [generalized quantum theory](#)

details

Bistable perception of ambiguous stimuli: the Necker cube



spontaneous switches between two possible 3-D representations  
at a time scale of some seconds

## Misra and Sudarshan (1977): Quantum Zeno Effect

- Two kinds of processes in an unstable two-state system:

“observation”:  $\sigma_3 = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}$  switching dynamics:  $\sigma_1 = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}$

## Misra and Sudarshan (1977): Quantum Zeno Effect

- Two kinds of processes in an unstable two-state system:

“observation”:  $\sigma_3 = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}$  switching dynamics:  $\sigma_1 = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}$

$$\sigma_1 \sigma_3 \neq \sigma_3 \sigma_1$$

## Misra and Sudarshan (1977): Quantum Zeno Effect

- Two kinds of processes in an unstable two-state system:

“observation”:  $\sigma_3 = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}$  switching dynamics:  $\sigma_1 = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}$

$$\sigma_1 \sigma_3 \neq \sigma_3 \sigma_1$$

- The switching dynamics is a continuous rotation according to

$$U(t) = e^{iHt} = \begin{pmatrix} \cos gt & i \sin gt \\ i \sin gt & \cos gt \end{pmatrix},$$

with  $H = g\sigma_1$ , and  $t_o = 1/g$  characterizes the decay time of the system.

## Misra and Sudarshan (1977): Quantum Zeno Effect

- Two kinds of processes in an unstable two-state system:

“observation”:  $\sigma_3 = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}$  switching dynamics:  $\sigma_1 = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}$

$$\sigma_1 \sigma_3 \neq \sigma_3 \sigma_1$$

- The switching dynamics is a continuous rotation according to

$$U(t) = e^{iHt} = \begin{pmatrix} \cos gt & i \sin gt \\ i \sin gt & \cos gt \end{pmatrix},$$

with  $H = g\sigma_1$ , and  $t_o = 1/g$  characterizes the decay time of the system.

- “Observation” process is a projection  $P_+$  or  $P_-$  onto one of the two eigenstates of  $\sigma_3$ .

- System dynamics without observations: Probability that the system is in state  $|+\rangle$  at time  $t$  if it was in  $|+\rangle$  at  $t = 0$ :

$$w_1(t) = \cos^2(gt)$$

- System dynamics without observations: Probability that the system is in state  $|+\rangle$  at time  $t$  if it was in  $|+\rangle$  at  $t = 0$ :

$$w_1(t) = \cos^2(gt)$$

- Successive observations at intervals  $\Delta T$ : Probability that the system is in state  $|+\rangle$  at time  $t = N \cdot \Delta T$  if it was in  $|+\rangle$  at  $t = 0$ :

$$\begin{aligned} w_N(t) &= (\cos^2(g\Delta T))^N \\ &\approx \exp(-g^2\Delta T^2 \cdot N) = \exp\left(-\frac{\Delta T}{t_0^2} t\right) \end{aligned}$$



- System dynamics without observations: Probability that the system is in state  $|+\rangle$  at time  $t$  if it was in  $|+\rangle$  at  $t = 0$ :

$$w_1(t) = \cos^2(gt)$$

- Successive observations at intervals  $\Delta T$ : Probability that the system is in state  $|+\rangle$  at time  $t = N \cdot \Delta T$  if it was in  $|+\rangle$  at  $t = 0$ :

$$\begin{aligned}w_N(t) &= (\cos^2(g\Delta T))^N \\ &\approx \exp(-g^2\Delta T^2 \cdot N) = \exp(-\frac{\Delta T}{t_0^2} t)\end{aligned}$$

- Effect of observations: stabilization of the system in its unstable states, “dwell time” increases from unperturbed  $t_0$  to an average time  $\langle T \rangle$ :

$$\langle T \rangle \approx t_0^2 / \Delta T$$

## From Quantum Zeno to Necker-Zeno

- States  $|+\rangle$  and  $|-\rangle$  correspond to the cognitive states in the two possible representations of the Necker cube.

## From Quantum Zeno to Necker-Zeno

- States  $|+\rangle$  and  $|-\rangle$  correspond to the cognitive states in the two possible representations of the Necker cube.
- Two complementary processes:
  - (i) unperturbed switching dynamics with characteristic time  $t_0$ ,
  - (ii) projection into a representation due to successive “updates” ( $\Delta T$ ).

## From Quantum Zeno to Necker-Zeno

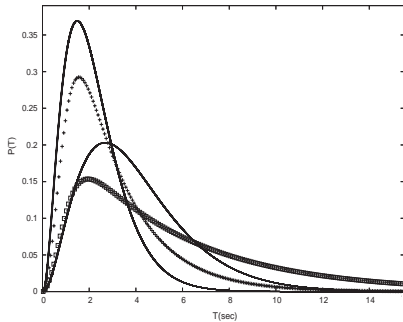
- States  $|+\rangle$  and  $|-\rangle$  correspond to the cognitive states in the two possible representations of the Necker cube.
- Two complementary processes:
  - (i) unperturbed switching dynamics with characteristic time  $t_0$ ,
  - (ii) projection into a representation due to successive “updates” ( $\Delta T$ ).
- Associated cognitive time scales:
  - intrinsic update interval  $\Delta T \approx 30$  msec  
 (sequentialization of successive stimuli, wagon wheel illusion)
  - $t_0 \approx 300$  msec (time for a stimulus to become conscious, P300)
  - $\langle T \rangle \approx 3$  sec (average “dwell time” for bistable states / representations)

Observed  $\Gamma$ -distribution of dwell times  $T$ :  $P(T) \propto T^b \exp(-\gamma T)$

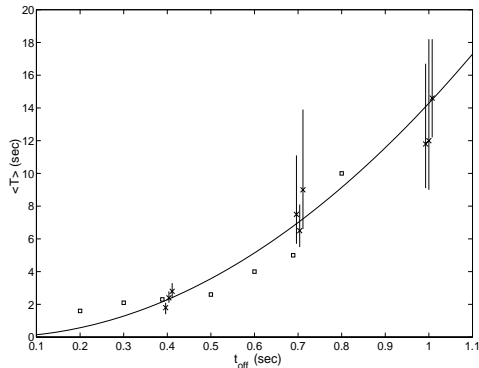
Model so far has  $b = 0$  (purely exponential decay of  $P(T)$ ),  
refine with initial behavior due to effects of attention:

(a) increasing  $\Delta T$ , (b) decreasing  $t_o$ .

- solid lines:  $\Gamma$ -distribution with  $b = 2$  and  $t_o = 300$  msec for  $\Delta T = 70$  msec (highest maximum) and  $\Delta T = 30$  msec
- $P(T)$  according to Necker-Zeno model with decreasing  $t_o$  for  $\Delta T = 70$  msec (crosses) and  $\Delta T = 30$  msec (squares)



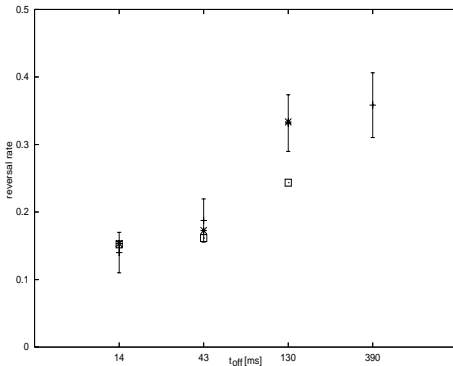
Dwell times  $\langle T \rangle$  for off-times  $t_{\text{off}} > 300$  msec



experimental values: crosses  
from Kornmeier and Bach  
(2004), squares from Orbach  
*et al.* (1966)

plotted curve according to  
the Necker-Zeno model for  
 $\Delta T \approx 70$  msec

Reversal rates  $1/\langle T \rangle$  for off-times  $t_{\text{off}} < 300$  msec



experimental values with  
error bars: from Kornmeier  
*et al.* (2007)

asterisks: best fit according  
to the Necker-Zeno model,  
yielding  $\Delta T \approx 16$  msec  
and  $t_0 \approx 210$  msec

squares: values for  $\Delta T = 30$   
msec and  $t_0 = 300$  msec

- Question: Can mental events always be uniquely assigned to instances without temporal extension?



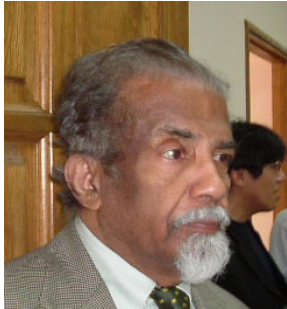
- Question: Can mental events always be uniquely assigned to instances without temporal extension?
- Bergson, James, Whitehead, etc., specious present, actual occasion, etc.: temporally extended events within which no further temporal localization (or segmentation) is possible. → **temporal nonlocality**

- Question: Can mental events always be uniquely assigned to instances without temporal extension?
- Bergson, James, Whitehead, etc., specious present, actual occasion, etc.: temporally extended events within which no further temporal localization (or segmentation) is possible. → **temporal nonlocality**
- In quantum mechanics, nonlocality is implied by non-commutative operations and can be tested experimentally. Bell's inequalities assume locality so that their violation demonstrates nonlocality.

- Question: Can mental events always be uniquely assigned to instances without temporal extension?
- Bergson, James, Whitehead, etc., specious present, actual occasion, etc.: temporally extended events within which no further temporal localization (or segmentation) is possible. → **temporal nonlocality**
- In quantum mechanics, nonlocality is implied by non-commutative operations and can be tested experimentally. Bell's inequalities assume locality so that their violation demonstrates nonlocality.
- Violations of temporal Bell inequalities would indicate temporal nonlocality (but in quantum mechanics time and dynamics are commutative).

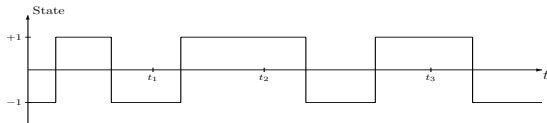
- Question: Can mental events always be uniquely assigned to instances without temporal extension?
- Bergson, James, Whitehead, etc., specious present, actual occasion, etc.: temporally extended events within which no further temporal localization (or segmentation) is possible. → **temporal nonlocality**
- In quantum mechanics, nonlocality is implied by non-commutative operations and can be tested experimentally. Bell's inequalities assume locality so that their violation demonstrates nonlocality.
- Violations of temporal Bell inequalities would indicate temporal nonlocality (but in quantum mechanics time and dynamics are commutative).
- In the Necker-Zeno model there are two kinds of non-commuting dynamics, so there is a chance to violate temporal Bell inequalities in bistable perception.

## Sudarshan (1983)

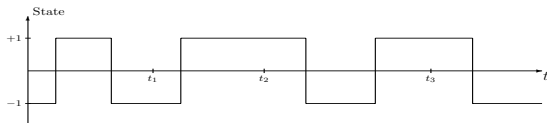


... a mode of awareness in which “sensations, feelings, and insights are not neatly categorized into chains of thoughts, nor is there a step-by-step development of a logical-legal argument-to-conclusion. Instead, patterns appear, interweave, coexist; and sequencing is made inoperative. Conclusion, premises, feelings, and insights coexist in a manner defying temporal order.”

- Specify three different instances  $t_1$ ,  $t_2$ ,  $t_3$  in a classical trajectory in which the state of the system at  $t_i$  is  $s(t_i) = \{+1, -1\}$ .

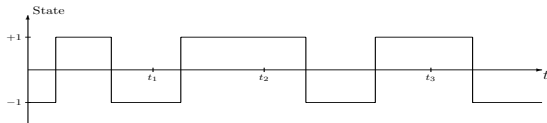


- Specify three different instances  $t_1$ ,  $t_2$ ,  $t_3$  in a classical trajectory in which the state of the system at  $t_i$  is  $s(t_i) = \{+1, -1\}$ .



- Any classical trajectory falls into one of  $2^3 = 8$  possible classes: 111, 11-1, 1-11, -111, 1-1-1, -11-1, -1-11, -1-1-1.

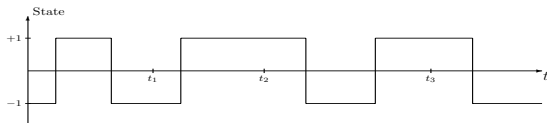
- Specify three different instances  $t_1$ ,  $t_2$ ,  $t_3$  in a classical trajectory in which the state of the system at  $t_i$  is  $s(t_i) = \{+1, -1\}$ .



- Any classical trajectory falls into one of  $2^3 = 8$  possible classes: 111, 11-1, 1-11, -111, 1-1-1, -11-1, -1-11, -1-1-1.
- Define  $N^-(t_i, t_j)$  as the number of cases with  $s(t_i) \neq s(t_j)$ , hence  $s(t_i)s(t_j) = -1$ , for each of the 8 possible trajectories.



- Specify three different instances  $t_1, t_2, t_3$  in a classical trajectory in which the state of the system at  $t_i$  is  $s(t_i) = \{+1, -1\}$ .



- Any classical trajectory falls into one of  $2^3 = 8$  possible classes: 111, 11-1, 1-11, -111, 1-1-1, -11-1, -1-11, -1-1-1.
- Define  $N^-(t_i, t_j)$  as the number of cases with  $s(t_i) \neq s(t_j)$ , hence  $s(t_i)s(t_j) = -1$ , for each of the 8 possible trajectories.
- For each trajectory,  $N^-(t_1, t_3) \leq N^-(t_1, t_2) + N^-(t_2, t_3)$ .  
Normalize  $N$  to  $p$ , replace  $(t_i, t_j)$  by  $(t_j - t_i)$ :

$$p^-(t_3 - t_1) \leq p^-(t_2 - t_1) + p^-(t_3 - t_2) \quad (\text{temporal Bell inequality})$$

- In the Necker-Zeno model, the probability for state  $|-\rangle$  at time  $t_2$  under the condition of state  $|+\rangle$  at time  $t_1$  (and vice versa) is:

$$w_{+-}(t_1, t_2) = w_{-+}(t_2, t_1) = \sin^2 g(t_2 - t_1)$$

- In the Necker-Zeno model, the probability for state  $|-\rangle$  at time  $t_2$  under the condition of state  $|+\rangle$  at time  $t_1$  (and vice versa) is:

$$w_{+-}(t_1, t_2) = w_{-+}(t_2, t_1) = \sin^2 g(t_2 - t_1)$$

- Then  $p^-(t_1, t_2)$  for anti-correlated states at  $t_1$  and  $t_2$  is:

$$p^-(t_1, t_2) = 1/2(w_{+-}(t_1, t_2) + w_{-+}(t_1, t_2)) = \sin^2 g(t_2 - t_1)$$

- In the Necker-Zeno model, the probability for state  $|-\rangle$  at time  $t_2$  under the condition of state  $|+\rangle$  at time  $t_1$  (and vice versa) is:

$$w_{+-}(t_1, t_2) = w_{-+}(t_2, t_1) = \sin^2 g(t_2 - t_1)$$

- Then  $p^-(t_1, t_2)$  for anti-correlated states at  $t_1$  and  $t_2$  is:

$$p^-(t_1, t_2) = 1/2(w_{+-}(t_1, t_2) + w_{-+}(t_1, t_2)) = \sin^2 g(t_2 - t_1)$$

- For  $\tau := t_3 - t_2 = t_2 - t_1$ , Bell's inequality turns into the sublinearity condition

$$p^-(2\tau) \leq 2p^-(\tau),$$

maximally violated for  $g\tau = \pi/6$  ( $\sin^2 g 2\tau = 3/4$ ,  $\sin^2 g \tau = 1/4$ ).

- In the Necker-Zeno model, the probability for state  $|-\rangle$  at time  $t_2$  under the condition of state  $|+\rangle$  at time  $t_1$  (and vice versa) is:

$$w_{+-}(t_1, t_2) = w_{-+}(t_2, t_1) = \sin^2 g(t_2 - t_1)$$

- Then  $p^-(t_1, t_2)$  for anti-correlated states at  $t_1$  and  $t_2$  is:

$$p^-(t_1, t_2) = 1/2(w_{+-}(t_1, t_2) + w_{-+}(t_1, t_2)) = \sin^2 g(t_2 - t_1)$$

- For  $\tau := t_3 - t_2 = t_2 - t_1$ , Bell's inequality turns into the sublinearity condition

$$p^-(2\tau) \leq 2p^-(\tau),$$

**maximally violated** for  $g\tau = \pi/6$  ( $\sin^2 g 2\tau = 3/4$ ,  $\sin^2 g \tau = 1/4$ ).

- For  $t_0 = 1/g \approx 300$  ms we obtain  $\tau = \pi/6 \cdot t_0 \approx 157$  ms as the optimal time difference between measurements of  $s(t_i)$ .

- In the Necker-Zeno model, the probability for state  $|-\rangle$  at time  $t_2$  under the condition of state  $|+\rangle$  at time  $t_1$  (and vice versa) is:

$$w_{+-}(t_1, t_2) = w_{-+}(t_2, t_1) = \sin^2 g(t_2 - t_1)$$

- Then  $p^-(t_1, t_2)$  for anti-correlated states at  $t_1$  and  $t_2$  is:

$$p^-(t_1, t_2) = 1/2(w_{+-}(t_1, t_2) + w_{-+}(t_1, t_2)) = \sin^2 g(t_2 - t_1)$$

- For  $\tau := t_3 - t_2 = t_2 - t_1$ , Bell's inequality turns into the sublinearity condition

$$p^-(2\tau) \leq 2p^-(\tau),$$

**maximally violated** for  $g\tau = \pi/6$  ( $\sin^2 g 2\tau = 3/4$ ,  $\sin^2 g \tau = 1/4$ ).

- For  $t_0 = 1/g \approx 300$  ms we obtain  $\tau = \pi/6 \cdot t_0 \approx 157$  ms as the optimal time difference between measurements of  $s(t_i)$ .
- Problem: measurements must be **as non-invasive as possible** to establish a significant violation of Bell's inequality.

H. Atmanspacher, H. Römer, H. Walach (2002): Weak quantum theory: Complementarity and entanglement in physics and beyond. *Foundations of Physics* **32**, 379–406.

H. Atmanspacher, T. Filk, H. Römer (2004): Quantum Zeno features of bistable perception. *Biological Cybernetics* **90**, 33–40.

H. Atmanspacher, M. Bach, T. Filk, J. Kornmeier, H. Römer (2008): Cognitive time scales in a Necker-Zeno model for bistable perception. *Open Cybernetics and Systemics Journal* **2**, 234–251.

H. Atmanspacher, T. Filk, H. Römer (2008): Complementarity in bistable perception. In *Recasting Reality. Wolfgang Pauli's Philosophical Ideas and Contemporary Science*, ed. by H. Atmanspacher and H. Primas, Springer, Berlin, pp. 135–150.

H. Atmanspacher, T. Filk (2010): A proposed test of temporal nonlocality in bistable perception. *Journal of Mathematical Psychology* **54**, 314–321.

B. Misra, E.C.G. Sudarshan (1977): The Zeno's paradox in quantum theory. *Journal of Mathematical Physics* **18**, 756–763.

Introduction

Necker-Zeno Model for Bistable Perception

Empirical Confirmation

Temporal Nonlocality

**Selected References**



- **Observables**  $A \in \mathcal{A}$  are (identifiable with) mappings  $A : \mathcal{Z} \mapsto \mathcal{Z}$  which associate to every **state**  $z \in \mathcal{Z}$  another state  $A(z)$ .
- To every observable  $A$  belongs a set  $specA$  of possible outcomes of an evaluation (e.g., “measurement”) of  $A$ .
- With  $A$  and  $B$ , also  $A \circ B$  is an observable.  
(An addition of observables is not defined.)
- There is a unit observable  $\mathbb{1}$ ,  $spec\mathbb{1} = \{\text{true}\}$ , such that:  
 $\mathbb{1}A = A\mathbb{1} \quad \forall A \in \mathcal{A}$ .
- For a zero state  $o$  and a zero observable  $O$ ,  $specO = \{\text{false}\}$ , we have:  
 $A(o) = o, \quad AO = OA = O, \quad \forall A \in \mathcal{A},$   
 $O(z) = o \quad \forall z \in \mathcal{Z}$ .
- Observables  $P$  with  $specP = \{\text{true}, \text{false}\}$  are **propositions** with the operations negation, conjunction, adjunction as usual.

- $\mathcal{A}$  has the structure of a monoid, generally non-commutative. The non-commutative case implies the concepts of:  
complementarity (incompatibility) of observables, dispersive states;  
entanglement (holistic correlations) among observables.  
(Cf. partially Boolean algebra of propositions.)

- $\mathcal{A}$  has the structure of a monoid, generally non-commutative. The non-commutative case implies the concepts of:  
complementarity (incompatibility) of observables, dispersive states;  
entanglement (holistic correlations) among observables.  
(Cf. partially Boolean algebra of propositions.)
- Generalized QT provides room for both ontic and epistemic interpretations. An ontic interpretation of complementarity and entanglement arises if pure states associated with incompatible observables are not dispersion-free.

- $\mathcal{A}$  has the structure of a monoid, generally non-commutative. The non-commutative case implies the concepts of:  
complementarity (incompatibility) of observables, dispersive states;  
entanglement (holistic correlations) among observables.  
(Cf. partially Boolean algebra of propositions.)
- Generalized QT provides room for both ontic and epistemic interpretations. An ontic interpretation of complementarity and entanglement arises if pure states associated with incompatible observables are not dispersion-free.
- The axiomatic framework of generalized QT does not prescribe the decomposition of a system  $\Sigma$  into subsystems. In particular there is no tensor product construction for composite systems.

- $\mathcal{A}$  has the structure of a monoid, generally non-commutative. The non-commutative case implies the concepts of:  
complementarity (incompatibility) of observables, dispersive states;  
entanglement (holistic correlations) among observables.  
(Cf. partially Boolean algebra of propositions.)
- Generalized QT provides room for both ontic and epistemic interpretations. An ontic interpretation of complementarity and entanglement arises if pure states associated with incompatible observables are not dispersion-free.
- The axiomatic framework of generalized QT does not prescribe the decomposition of a system  $\Sigma$  into subsystems. In particular there is no tensor product construction for composite systems.
- For the dynamical evolution of  $\Sigma$  one may assume a one-parameter (semi-) group of endomorphisms. However, there is no prescribed kind of dynamical evolution for subsystems of  $\Sigma$  and their interaction.

“Most generalized” QT does not use key features of ordinary QT



no algebra,



no space,



no rule,



no action,



no uncertainty,



no equation,



no inequalities

[← return](#)